Tensor network studies of SrCu₂(BO₃)₂ under pressure and in a magnetic field

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Shastry-Sutherland model







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Outline

- Part I: SCBO under pressure at finite temperature
 - ✦ Critical point at finite temperature, analogous to the critical point of water

J. L. Jiménez, S. P. G. Crone, E. Fogh, M. E. Zayed, R. Lortz, E. Pomjakushina,K. Conder, A. M. Läuchli, L. Weber, S. Wessel, A. Honecker, B. Normand, C. Rüegg,PC, H. M. Rønnow, and F. Mila, Nature 592, 370 (2021)

Part II: SCBO under extreme conditions of field & pressure

New type of 1/5 plateau and supersolid phases in the (p, H) phase diagram
 Z. Shi, S. Dissanayake, PC, W. Steinhardt, D. Graf, D. M. Silevitch, H. A. Dabkowska,
 T. F. Rosenbaum, F. Mila, S. Haravifard, Nat Commun 13, 1 (2022)

Part III: SCBO up to the saturation magnetic field

Close agreement between iPEPS and experiments & insights into ultrasound velocity
 T. Nomura, PC, A. Miyata, S. Zherlitsyn, Y. Ishii, Y. Kohama, Y. Matsuda, A. Ikeda,
 C. Zhong, H. Kageyama, F. Mila, arXiv:2209.07652 (to appear in Nat. Comm.)

Part IV: Role of interlayer coupling in SCBO

Phase diagram of the SSM with interlayer coupling with new iPEPS approach
 P. Vlaar, PC, PRL 130, 130601 (2023); arxiv:2302.07894

The Shastry-Sutherland model and SrCu₂(BO₃)₂



The Shastry-Sutherland model and SrCu₂(BO₃)₂



Magnetization plateaus

SrCu₂(BO₃)₂ in a magnetic field exhibits several magnetization plateaus



The SSM has almost localized triplet excitations [Miyahara&Ueda'99, Kageyama et al. '00]

Triplets repel each other (on the mean-field level)

Common assumption:

magnetization plateaus correspond to crystals of localized triplets!



Magnetization plateaus below the 1/4 plateau



★ Crystals of triplet-bound states PC, F. Mila, PRL 112 (2014)

Many experimental / theoretical studies

Kageyama et al, PRL 82 (1999) Onizuka et al, JPSJ 69 (2000) Kageyama et al, PRL **84** (2000) Kodama et al, Science **298** (2002) Takigawa et al, Physica 27 (2004) Levy et al, EPL 81 (2008) Sebastian et al, PNAS 105 (2008) Isaev et al, PRL 103 (2009) Jaime et al, PNAS **109** (2012) Takigawa et al, PRL **110** (2013) Matsuda et al, PRL 111 (2013) Miyahara and K. Ueda, PRL 82 (1999) Momoi and Totsuka, PRB 61 (2000) Momoi and Totsuka, PRB 62 (2000) Fukumoto and Oguchi, JPSJ 69 (2000) Fukumoto, JPSJ 70 (2001) Miyahara and Ueda, JPCM 15 (2003) Miyahara, Becca and Mila, PRB 68 (2003) Dorier, Schmidt, and Mila, PRL 101 (2008) Abendschein & Capponi, PRL 101 (2008) Takigawa et al, JPSJ 79 (2010) Nemec et al, PRB 86 (2012) Matsuda et al., PRL 111 (2013)

SrCu₂(BO₃)₂ under pressure



Bettler, et al., Phys. Rev. Research 2, 012010 (2020) Cui et al., arxiv:2204.08133



MPS & PEPS

D

MPS

Matrix-product state



PEPS (TPS)

projected entangled-pair state (tensor product state)



Physical indices (lattices sites)

S. R. White, PRL 69, 2863 (1992) Fannes et al., CMP 144, 443 (1992) Östlund, Rommer, PRL 75, 3537 (1995) F. Verstraete, J. I. Cirac, cond-mat/0407066 Nishio, Maeshima, Gendiar, Nishino, cond-mat/0401115

Infinite PEPS (iPEPS)

D iMPS

infinite matrix-product state



2D

iPEPS

infinite projected entangled-pair state



Jordan, Orus, Vidal, Verstraete, Cirac, PRL (2008)

Work directly in the thermodynamic limit:
 No finite size and boundary effects!

Infinite PEPS (iPEPS)

iMPS

infinite matrix-product state



2D

iPEPS

infinite projected entangled-pair state



Jordan, Orus, Vidal, Verstraete, Cirac, PRL (2008)

Work directly in the thermodynamic limit:
 No finite size and boundary effects!

iPEPS with arbitrary unit cells

ID iMPS

infinite matrix-product state







iPEPS

with arbitrary unit cell of tensors



here: 4x2 unit cell

PC, White, Vidal, Troyer, PRB 84 (2011)

★ Run simulations with different unit cell sizes and compare variational energies

iPEPS ground state simulations

• Many applications to challenging problems, including frustrated spin, SU(N), bosonic systems, t-J / Hubbard models, and more, see e.g.:

- P. Corboz and F. Mila, PRB 87 (2013); PRL 112 (2014)
- Z.-C. Gu, H.-C. Jiang, D. N. Sheng, H. Yao, L. Balents and X.-G. Wen, PRB 88 (2013)
- J. Osorio Iregui, P. Corboz and M. Troyer, PRB 90 (2014)
- P. Corboz, T. Rice and M. Troyer, PRL 113 (2014)
- T. Picot and D. Poilblanc, PRB 91 (2015)
- T. Picot, M. Ziegler, R. Orús and D. Poilblanc, PRB 93 (2016)
- P. Nataf, M. Lajkó, P. Corboz, A. M. Läuchli, K. Penc and F. Mila, PRB 93 (2016)
- H. Liao, Z. Xie, J. Chen, Z. Liu, H. Xie, R. Huang, B. Normand and T. Xiang, PRL 118 (2017)
- B.-X. Zheng, et al., Science 358, 1155 (2017)
- I. Niesen and P. Corboz, PRB 95 (2017); SciPost Physics 3, 030 (2017); Rev. B 97, 245146 (2018)
- R. Haghshenas, W.-W. Lan, S.-S. Gong, and D. N. Sheng, PRB 97 (2018)
- J.-Y. Chen, L. Vanderstraeten, S. Capponi, and D. Poilblanc, PRB 98 (2018)
- S. S. Jahromi and R. Orús, PRB 98 (2018)
- H.-Y. Lee and N. Kawashima, PRB 97 (2018)
- H. Yamaguchi, Y. Sasaki, T. Okubo, et al., PRB 98, 094402 (2018)
- R. Haghshenas, S.-S. Gong, and D. N. Sheng, PRB 99, 174423 (2019)
- S. S. Chung and P. Corboz, PRB 100 (2019)
- B. Ponsioen, S. S. Chung, and P. Corboz, PRB 100 (2019)
- C. Boos, S. P. G. Crone, I. A. Niesen, P. Corboz, K. P. Schmidt, and F. Mila, PRB 100 (2019)
- Z. Shi, et al, Nature Communications 10, 2439 (2019)
- A. Kshetrimayum, C. Balz, B. Lake, and J. Eisert, ArXiv:1904.00028 (2019)
- H.-Y. Lee, R. Kaneko, T. Okubo, and N. Kawashima, PRL 123, 087203 (2019)
- O. Gauthé, S. Capponi, M. Mambrini, and D. Poilblanc, PRB 101, 205144 (2020)
- H.-Y. Lee, R. Kaneko, L. E. Chern, T. Okubo, Y. Yamaji, N. Kawashima, and Y. B. Kim, Nature Communications 11 (2020)
- W.-Y. Liu, S.-S. Gong, Y.-B. Li, D. Poilblanc, W.-Q. Chen, and Z.-C. Gu, ArXiv:2009.01821 (2020)
- J.-Y. Chen, S. Capponi, A. Wietek, M. Mambrini, N. Schuch, and D. Poilblanc, PRL 125, 017201 (2020)
- J. Hasik, D. Poilblanc, and F. Becca, SciPost Physics 10, 012 (2021)
- ... and many more ...

S. Dusuel, M. Kamfor, R. Orús, K. P. Schmidt, and J. Vidal, PRL 106, 107203 (2011)

P. Corboz, A. M. Läuchli, K. Penc, M. Troyer and F. Mila, PRL 107 (2011)

H. H. Zhao, C. Xu, Q. N. Chen, Z. C. Wei, M. P. Qin, G. M. Zhang and T. Xiang, PRB 85 (2012)

P. Corboz, M. Lajkó, A. M. Läuchli, K. Penc and F. Mila, PRX 2 (2012)

Finite temperature simulations with iPEPS

Methodological developments (2D):

Li et al. PRL 106 (2011); Czarnik et al. PRB 86 (2012); Czarnik & Dziarmaga PRB 90 (2014); Czarnik & Dziarmaga PRB 92 (2015); Czarnik et al. PRB 94 (2016); Dai et al PRB 95 (2017); Kshetrimayum, Rizzi, Eisert, Orus, PRL 122 (2019), P. Czarnik, J. Dziarmaga, PC, PRB 99 (2019), ...



 $\hat{
ho}(eta) = \hat{
ho}^{\dagger}(eta)$ by construction

Finite temperature simulation benchmarks

Wietek, PC, Wessel, Normand, Mila, and Honecker, PRR 1 (2019)

- Benchmarks in the dimer phase of the Shastry-Sutherland model
- Comparison between ED, TPQ, QMC, iPEPS



Specific heat: experiments vs iPEPS





Jiménez, Crone, et al., Nature 592, 370 (2021)

Correlation length



Jump in $<S \cdot S >$ on dimer



Clear evidence of a first order line with a critical point compatible with the 2D Ising universality class

2D classical Ising model in a field



Open challenges



🔷 🔿 🛧 Dynamics Order 🛛 🗖 200 80 50 g 60 Energy (K) rature full plaquette 100 📆 40 phase (FPP) 50 20 AFM Plaquette Dimer AFM 0 0 10 20 30 40 50 60 Pressure (kbar)

too low T to obtain reliable data with iPEPS (currently)

NMR & INS experiments: full plaquette phase (FPP), not empty plaquette phase (EPP) Zayed, et al., Nat. Phys. 13, 962 (2017) Vaki, et al. JPSJ 76 (2007) Cui et al., arxiv:2204.08133



Competing plaquette phases

Distorted Shastry-Sutherland model: competition between EPP and FPP phase





Small deformation leads to FPP phase!

But precise model still unclear...

Part II: SCBO under extreme conditions of field & pressure

Shi, Dissanayake, PC, William Steinhardt, Graf, Silevitch, Dabkowska, Rosenbaum, Mila, Haravifard, Nat Commun 13, 1 (2022)

 Experiments: tunnel diode oscillator (TDO) technique

$$\frac{df}{dH} \propto -\frac{d^2 M}{dH^2}$$

- Non-zero df/dH ↔ slope change in M
- Identify anomalies \rightarrow phase transitions
- Compare with iPEPS phase diagram



SrCu₂(BO₃)₂ under extreme conditions of field & pressure

Shi, Dissanayake, PC, William Steinhardt, Graf, Silevitch, Dabkowska, Rosenbaum, Mila, Haravifard, Nat Commun 13, 1 (2022)



Nature of the 1/5 plateau



- Vertical stripes separated by dimer singlets (along red dashed lines)
- Strong triplets in the center of each stripe, with a weaker pair of triplets in between
 - Neither a crystal of triplets, nor a crystal of triplet bound states!
- Full-plaquette formation: reminiscent of full plaquette phase (FPP)
 - ✦ FPP: triplets on dimers within plaquette and singlets on dimers outside of plaquette, where the triplets form effective S=1 Haldane chains
- Effective description: S=1 diamond chain with m=2/3
- Also found in a thin SSM tube made of 2 orthogonal dimer chains

Manmana, Picon, Schmidt, and Mila, EPL 94, 67004 (2011)

Nature of the 10x2 supersolid



- Descendant of the I/5 plateau state
- Alternating rotation of the spins of successive stripes clockwise or counterclockwise by 90 degrees
- Finite component in the field direction, also on the boundary between stripes

Full plaquette physics appearing at finite magnetic field!

Part III: SrCu₂(BO₃)₂ up to the saturation field

Nomura, PC, Miyata, Zherlitsyn, Ishii, Kohama, Matsuda, Ikeda, Zhong, Kageyama, Mila, arXiv:2209.07652 (to appear in Nat. Comm.)

- Experiments: ultrahigh fields up to 150T!
- Ultrasound velocity & magnetostriction
- Identified new anomalies above the I/2 plateau at II6T, I27T, and I39T
- Saturation field: I 39T

- Comparison with iPEPS results
- Strong decrease of the sound velocity of the c₆₆ acoustic in the 1/2 plateau?



iPEPS phase diagram













Results for J'/J=0.63



Anomalies H_{c6}, H_{c8}, and H_{c9} compatible with iPEPS results

Reduction of ultrasound velocity in 1/2 plateau?



Strain of c66 mode:



- Strong reduction in 1/2 plateau
- Tiny reduction in 1/3 plateau

- Both E' and E'' of the magnetic energy contribute to the elastic constant
- E' has large magnitude in 1/2, but vanishing for 1/3 plateau



Checkerboard structure with positive (triplet) and negative (singlet) bonds → contributions to E' add up → large magnitude



odd periodicity \rightarrow cancellation

Reduction of ultrasound velocity in 1/2 plateau?



Strain of c66 mode:



- Strong reduction in 1/2 plateau
- Tiny reduction in 1/3 plateau

Both E' and E'' of the magnetic energy contribute to the elastic constant

Fit:
$$\lambda_1 E' + \lambda_2 E'' \sim \Delta c/c$$
, $c = \rho v^2$

Plateau	E'	$E^{\prime\prime}$	$\Delta v/v_0$	$\Delta v/v_0$
			iPEPS	Exp.
1/8	-0.014	-0.44	-0.13	-0.11(4)
1/4	-0.051	-0.31	-0.16	-0.19(3)
1/3	0	-0.18	-0.04	-0.03(2)
1/2	-0.23	-0.034	-0.49	-0.48(7)

Estimates in good agreement with experiment

Part IV: SSM with interlayer coupling

• Extent of the plaquette phase is smaller in experiments than in theory



Pressure model: Shi, et al. (2022)

 $p_0 \leftrightarrow J'/J = 0.63$ $p_c = 1.8GPa \leftrightarrow J'/J = 0.675$ J = 81.5KJ'(p), J(p): linear functions. J'(p) changes by 5% between p_0 and p_c (ESR) Sakurai, et al., J. Phys. Soc. Jpn. 87, 033701 (2018)

SSM with interlayer coupling



AF phase becomes favored over plaquette phase with increasing J"

Predicted values for J" (no consensus yet)

✓ J"/J = 0.09 ... 0.21 from fits to susceptibility [Miyahara & Ueda (2000), Knetter et al (2000)]
 ✓ J"/J < 0.03 from ab-initio calculations [Radtke et al., PNAS 112 (2015)]

iPEPS for layered systems

Vlaar, PC, PRL 130 (2023)



Ansatz:

- 3D tensor network ansatz (coupled iPEPS)
- $D_{xy} > D_z$ for weak interlayer coupling
- $D_z = I \rightarrow \text{product state of iPEPSs}$



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Contraction:

- $D_z = I$: contract individual layers (2D)
- $D_z > I$: perform effective decoupling away from center $\rightarrow 2D$ contraction
- Interlayer correlations beyond meanfield level are included by the D_z > I bonds in the center
- Layered corner transfer matrix (LCTM) method



Benchmarks for 3D anisotropic Heisenberg model

Vlaar, PC, PRL 130 (2023)



- Substantial improvement from $D_z = 1$ to $D_z = 2$
- Values close to the extrapolated QMC result
- In agreement with more expensive full 3D contractions

Vlaar & PC, PRB 103, 205137 (2021)

Phase diagram: SSM with interlayer coupling



Phase diagram: SSM with interlayer coupling

Vlaar, PC, arxiv:2302.07894



Estimate for the strength of interlayer coupling: $J''/J \approx 0.03$

LCTM: promising approach also for other layered systems

Conclusion

- \checkmark SrCu₂(BO₃)₂ under pressure / in a magnetic field exhibits very rich physics!
- Finite temperature Ising critical point, analogous to critical point of water
- New type of 1/5 plateau and supersolid phases at high pressure & field \checkmark
- \checkmark Results up to saturation in good agreement with experiments & understanding of the reduction of ultrasound velocity in the 1/2 plateau
- \checkmark Reduction of plaquette phase due to interlayer coupling
- Progress with iPEPS: versatile tool for ground state and finite temperature calculations + extensions to layered systems

Thank you for your attention!

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